

REVIEW AND SUMMARY FOR CHAPTER 8

This chapter was devoted to the study of *dry friction*, that is, to problems involving rigid bodies which are in contact along *nonlubricated surfaces*.

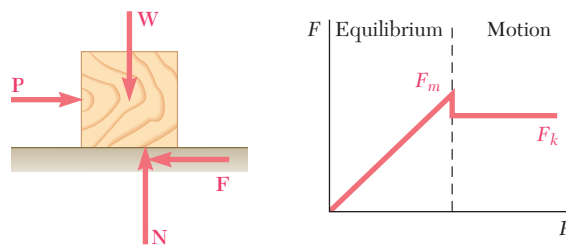


Fig. 8.16

Applying a horizontal force \mathbf{P} to a block resting on a horizontal surface [Sec. 8.2], we note that the block at first does not move. This shows that a *friction force* \mathbf{F} must have developed to balance \mathbf{P} (Fig. 8.16). As the magnitude of \mathbf{P} is increased, the magnitude of \mathbf{F} also increases until it reaches a maximum value F_m . If \mathbf{P} is further increased, the block starts sliding and the magnitude of \mathbf{F} drops from F_m to a lower value F_k . Experimental evidence shows that F_m and F_k are proportional to the normal component N of the reaction of the surface. We have

$$F_m = \mu_s N \quad F_k = \mu_k N \quad (8.1, 8.2)$$

where μ_s and μ_k are called, respectively, the *coefficient of static friction* and the *coefficient of kinetic friction*. These coefficients depend on the nature and the condition of the surfaces in contact. Approximate values of the coefficients of static friction were given in Table 8.1.

It is sometimes convenient to replace the normal force \mathbf{N} and the friction force \mathbf{F} by their resultant \mathbf{R} (Fig. 8.17). As the friction force increases and reaches its maximum value $F_m = \mu_s N$, the angle ϕ that \mathbf{R} forms with the normal to the surface increases and reaches a maximum value ϕ_s , called the *angle of static friction*. If motion actually takes place, the magnitude of \mathbf{F} drops to F_k ; similarly the angle ϕ drops to a lower value ϕ_k , called the *angle of kinetic friction*. As shown in Sec. 8.3, we have

$$\tan \phi_s = \mu_s \quad \tan \phi_k = \mu_k \quad (8.3, 8.4)$$

Static and kinetic friction

Angles of friction

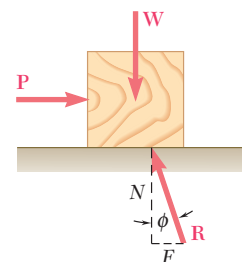


Fig. 8.17

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Problems involving friction

When solving equilibrium problems involving friction, we should keep in mind that the magnitude F of the friction force is equal to $F_m = \mu_s N$ *only if the body is about to slide* [Sec. 8.4]. *If motion is not impending, F and N should be considered as independent unknowns* to be determined from the equilibrium equa-

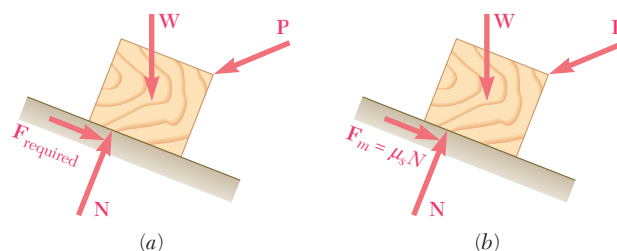


Fig. 8.18

tions (Fig. 8.18a). We should also check that the value of F required to maintain equilibrium is not larger than F_m ; if it were, the body would move and the magnitude of the friction force would be $F_k = \mu_k N$ [Sample Prob. 8.1]. On the other hand, *if motion is known to be impending, F has reached its maximum value $F_m = \mu_s N$* (Fig. 8.18b), and this expression may be substituted for F in the equilibrium equations [Sample Prob. 8.3]. When only three forces are involved in a free-body diagram, including the reaction \mathbf{R} of the surface in contact with the body, it is usually more convenient to solve the problem by drawing a force triangle [Sample Prob. 8.2].

When a problem involves the analysis of the forces exerted on each other by *two bodies A and B*, it is important to show the friction forces with their correct sense. The correct sense for the friction force exerted by B on A , for instance, is opposite to that of the *relative motion* (or impending motion) of A with respect to B [Fig. 8.6].

Wedges and screws

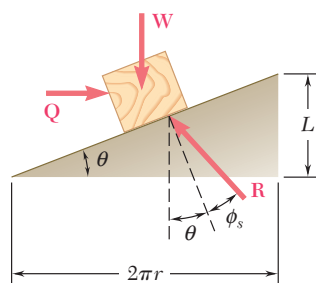


Fig. 8.19

In the second part of the chapter we considered a number of specific engineering applications where dry friction plays an important role. In the case of *wedges*, which are simple machines used to raise heavy loads [Sec. 8.5], two or more free-body diagrams were drawn and care was taken to show each friction force with its correct sense [Sample Prob. 8.4]. The analysis of *square-threaded screws*, which are frequently used in jacks, presses, and other mechanisms, was reduced to the analysis of a block sliding on an incline by unwrapping the thread of the screw and showing it as a straight line [Sec. 8.6]. This is done again in Fig. 8.19, where r denotes the *mean radius* of the thread, L is the *lead* of the screw, that is, the distance through which the screw advances in one turn, \mathbf{W} is the load, and Qr is equal to the torque exerted on the screw. It was noted that in the case of multiple-threaded screws the lead L of the screw is *not* equal to its pitch, which is the distance measured between two consecutive threads.

Other engineering applications considered in this chapter were *journal bearings* and *axle friction* [Sec. 8.7], *thrust bearings* and *disk friction* [Sec. 8.8], *wheel friction* and *rolling resistance* [Sec. 8.9], and *belt friction* [Sec. 8.10].

In solving a problem involving a *flat belt* passing over a fixed cylinder, it is important to first determine the direction in which the belt slips or is about to slip. If the drum is rotating, the motion or impending motion of the belt should be determined *relative* to the rotating drum. For instance, if the belt shown in Fig. 8.20 is

Belt friction

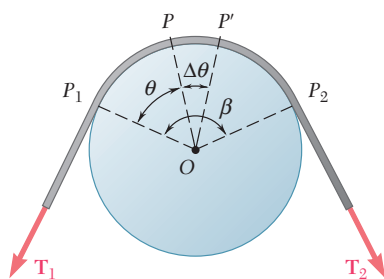


Fig. 8.20

about to slip to the right relative to the drum, the friction forces exerted by the drum on the belt will be directed to the left and the tension will be larger in the right-hand portion of the belt than in the left-hand portion. Denoting the larger tension by T_2 , the smaller tension by T_1 , the coefficient of static friction by μ_s , and the angle (in radians) subtended by the belt by β , we derived in Sec. 8.10 the formulas

$$\ln \frac{T_2}{T_1} = \mu_s \beta \quad (8.13)$$

$$\frac{T_2}{T_1} = e^{\mu_s \beta} \quad (8.14)$$

which were used in solving Sample Probs. 8.7 and 8.8. If the belt actually slips on the drum, the coefficient of static friction μ_s should be replaced by the coefficient of kinetic friction μ_k in both of these formulas.